The turbulence between geomechanics and vibroacoustics: what is common point?

S. Bonelli and P.-O. Mattei

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Fabien of course

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Three main parts.

Each corresponds to a work that was made possible under Fabien's impetus

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3. Singing risers (IRPHE-LMA)

Part 1 Erosion of a cohesive soil by turbulent flow Stéphane BONELLI

Part 2

Acoustic emission of a turbulent flow in a soil

Stéphane $\operatorname{BONELLI}$ and Pierre-Olivier MATTEI

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Simplified model for a soil experimental model - PhD thesis Ch. Jeanniot

Let's consider a cylindrical sample of soil of radius R_e and length L_e , pierced in its center by a cylindrical hole of radius R_i . This sample is embedded in a housing of radius R_p and length L_p .



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Can we measure at the surface of the housing the sound pressure created by the inner flow ?

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Can we measure at the surface of the housing the sound pressure created by the inner flow ?

is it possible to quantify the inner flow/inner hole ?

The acoustic pressure p(M) inside the soil is solution of

$$\begin{bmatrix} \Delta - i\eta \frac{\omega}{c_0^2} + \frac{\omega^2}{c_0^2} \end{bmatrix} p(M, \omega) = S(M', \omega), M, M' \in]0, R_e] \times]0, L_p]$$

$$p(M, \omega) = 0, x = 0, x = L_p \text{ soft boundary}$$

$$\partial_n p(M, \omega) = 0, r = R_p \text{ rigid boundary}$$

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Because of the simplified geometry, P(M) can be expressed as a sum of normal modes $\Psi_{mns}(M) = A_{mns} \exp(im\theta) \cos(n\pi x/L_p) J_m(\kappa_{ms}r)$ where $J'_m(\kappa_{ms}R_p) = 0$:

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As $\kappa_{0,0} = 0$, the first two resonances $f_{mns} = \omega_{mns}/(2\pi)$ of this system are given by $f_{0n0} = c_0 \times n/(2L_p)$, n = 1, 2, with $c_0 = 1500$ m/s and $L_p = m$: $f_{010} \approx 2500$ Hz and $f_{010} \approx 5000$ Hz

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Now, if the flow inside the soil sample is turbulent, it can be described by a Corcos model which power spectral density of the excitation is given by

$$\Phi_{\mathcal{S}}(Q,Q',\omega) = \Phi_{0}(\omega) \exp\left(\frac{-R_{i}|\theta - \theta'|}{L_{\theta}(\omega)}\right) \exp\left(\frac{-|x - x'|}{L_{x}(\omega)}\right) \exp\left(\frac{i\omega(x - x')}{U_{c}(\omega)}\right)$$

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The variable separation in both the Corcos model and pressure field allows analytical expression for the power spectral density of the pressure inside the soil sample given by linear combination of Bessel and trigonometric functions

$$\Phi_{\rho}(M, M'; \omega) = \Phi_{0}(\omega) \sum_{m=-\infty}^{\infty} I_{\theta m}(L_{\theta}, R_{i}) \\ \times \sum_{n, s, \nu, \sigma} \frac{J_{m}(\kappa_{ms}r)J_{m}(\kappa_{ms}R_{i})J_{m}(\kappa_{m\sigma}r')J_{m}(\kappa_{m\sigma}R_{i})}{A_{mns}A_{m\nu\sigma}(\omega^{2} - \imath\eta\omega - \omega_{mns}^{2})(\omega^{2} - \imath\eta\omega - \omega_{m\nu\sigma}^{2})} \\ \times \cos\left(\frac{n\pi x}{L_{\rho}}\right)\cos\left(\frac{\nu\pi x'}{L_{\rho}}\right)I_{n}^{\nu}(L_{x}, U_{c})\exp(\imath\theta - \theta')$$

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 $I_{\theta m}(L_{\theta}, R_i)$ and $I_{\mu}^{\nu}(L_x, U_c)$ are trigonometric functions (a few lines length...). By inverting the power spectral density of the pressure, it is possible (not so easy

indeed) to obtain U_c and R_i .

First experimental results



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Characteristic signal (left): above 1 kHz, Peak identification (right) by fitting a Lorenztian function.

First experimental results



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Characteristic signal (left): above 1 kHz, Peak identification (right) by fitting a Lorenztian function.

Peaks at 2340 Hz and 5470 Hz (close to those predicted)



Evolution of the amplitude of resonance peaks as a function of Reynolds Number

Evolution of the two resonances when Reynolds varies

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Model for the HET soil sample promising (but to be numerically tested)

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- New set of experiments to be started by the end of July

Model for the HET soil sample promising (but to be numerically tested)

- First experiments in line with the model
- New set of experiments to be started by the end of July
- Similar models for real dikes ?

Part 3 Singing risers

Pierre-Olivier MATTEI



Singing risers: introduction

• Under particular flow conditions, a corrugated pipe can "sing": applications in music (voice of the dragon) or in industry (singing risers/FLow Induced Pulsation).



- Noise on the platform : 110 dB
- Noise inside the pipe : 160 dB (vibration and fatigue issue for adjacent equipment)

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Flow rate limitation

1. FLIP mitigation using low frequency sound using "lab" experiment

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- origin of the whistling
- one or more noise sources
- at which localization(s)
- weak of strong aeroacoustic coupling ?
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"industrial riser" + high speed flow / high pressure + microphones + hot wire probe

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- influence of flow speed and static pressure on the FLIP onset
- one or more noise sources
- at which localization(s)
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FLIP mitigation using low frequency sound - Coll. U. Kristiansen (NTNU)



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Sound pressure recorded inside the pipe at a flow speed of about 20 m/s.



With mitigation at 14 Hz : signal 154 dB - audible signal 124 dB

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Time records of flow speed $SV_o(t)$ in m/s and acoustics sound pressure $SP_o(t)$ in Pa.

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• One subtracts to the signal (black) the LF component and its harmonics (red)



Fluctuations of flow speed and acoustics sound pressure.

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 \Rightarrow Very strong similarity of the residues



Gabor transform of flow speed and acoustics sound pressure

 \Rightarrow Shift of signals of about 0,007 s: flying time between entry and exit

 \Rightarrow Components separated by 0,014 s: round-trip flying time between entry and exit

Pearson's correlation between $SP_c(t)$ et $SV_c(t + \tau)$ vs τ calculated on 2^{17} samples for 6 000 time steps



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|r| = 0,5 at $\tau = 0,0069$.

The correlation fluctuates at 669 Hz.

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- ▶ the flow LF modulation limit the emission of pressure fluctuation
- ▶ its frequency can be tuned to limit amplification by reflection

FLIP measurement using industrial riser - PhD thesis G. Galeron

Experiments realized on a 18.34 m length "true" industrial riser (15.24 cm internal diameter) at CESAME test facility centre (Poitiers, France):

- ▶ Internal Pressure $p_i \in [1; 42]$ bars
- ► Flow velocities v₀ ∈ [5; 80] m/s
- 6 microphone probes Gras $p_i < 6$ bars
- 4 microphones Kulite $p_i < 50$ bars
- 2 Dantec Hot wires V₀ > 0 m/s
- $\blacktriangleright \Delta P$ probe, thermal probe, static pressure probe



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FLIP measurement using industrial riser - PhD thesis G. Galeron

Air supplying realized through emptying of a 200 m^3 tank at 200 bars ;



Fixture photograph - downstream coupler

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Example of time-frequency analysis around the maximum whistling of the riser (upstream)



Wavelet (Gabor) transform of an upstream microphone signal

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Fluctuations in time and amplitude. Upstream signal noisy (induced by fixture air supplying).

Example of time-frequency analysis around the maximum whistling of the riser (downstream)



Wavelet (Gabor) transform of a downstream microphone signal

Fluctuations in time, amplitude and frequency.

Downstream signal almost free from noise. The pipe acts as a pass-band filter around its transverse resonances.



Microphones 4, 5 and 6 between 449 s and 450 s. Strong whistling (about 170 dB !).

Wavelet (Gabor) transform of the microphone 4 signal

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Wavelet (Gabor) transform of the microphone 5 signal

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Between MG4 et MG5 ($d(MG_4; MG_5) = 0.4 \text{ m}$), we estimate a time shift of about $\Delta t = 5.7 \text{ ms}$, that is $L_{\Delta t}^{c_0} \approx 1.9 \text{ m}$ and $L_{\Delta t}^{V_0} \approx 0.39 \text{ m}$





Wavelet (Gabor) transform of the microphone 6 signal

Between MG4 et MG6 ($d(MG_5; MG_6) = 0.5 \text{ m}$), we estimate a time shift of about $\Delta t = 7 \text{ ms}$, that is $L_{\Delta t}^{c_0} \approx 2.3 \text{ m}$ and $L_{\Delta t}^{V_0} \approx 0.49 \text{ m}$

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Correlation analysis

Pearson's correlation of microphones 4 and 5 (t \in [450; 450] s) bandpass filtering between 800 Hz and 2000 Hz.

▶ 11 "oscillations" (ie 22 relative maxima) to reach the maximum correlation: $d(MG_4; MG_5) = 0.5 \text{ m} \Rightarrow 0.4/22 \approx 2 \text{ cm} \approx d_c$, correlation pitch

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The strong whistling inside the riser is dominated by the pseudo-noise which increase is induced by a retro-action of the internal acoustics field over the vortices caused by the corrugation

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- The strong whistling inside the riser is dominated by the pseudo-noise which increase is induced by a retro-action of the internal acoustics field over the vortices caused by the corrugation
- During whistling, the pseudo-noise spectrum is dominated by frequencies around that of the internal acoustics resonances

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Last but not least

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Many thanks to you, Fabien, for making possible all this work

